

a smaller chance of crossing the void gap and entering the adjacent non-void regions.

Therefore, one should undertake some numerical checks of the method described here before applying it to specific problems. I do not have an opportunity to do this, at present, and have submitted this letter trusting someone else will.

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Comments on "Generalizations of Fick's Law"*

In this letter we present several comments on a letter by Kostin¹. For brevity, we shall refer to "Eq. (n) of Ref. 1" as "(K-n)."

(I) The "generalized Fick's law" (K-11) is alleged to have the simplicity of Fick's law. We believe that this is not so, for the following reasons: 1) The quantity which appears under the derivative sign is $\bar{v}^2(z)N(z)$ and not $N(z)$. 2) The $\bar{k}_{tr}(z)$, which plays the role of a reciprocal diffusion coefficient, is not an intensive property of the medium^{2,3}. Even in a homogeneous medium, $\bar{k}_{tr}(z)$ may vary with z . 3) the $\bar{k}_{tr}(z)$ depends explicitly on the neutron distribution $n(z, v, \underline{\omega})$. In the time-dependent form (K-7), the $\bar{k}_{tr}(z, t)$ depends explicitly on $n(z, v, \underline{\omega}, t)$. It follows that $\bar{k}_{tr}(z, t)$ may be a function of time in a system with fixed composition and fixed boundaries.

(II) The definition (K-5) of $\alpha(z, v')$ is, with trivial changes in notation,

$$\int \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) v dv d\omega = v' \alpha(z, v') \Sigma_s(z, v').$$

This definition does not lead to (K-7) or (K-11); however a corrected definition

$$\int \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) v \mu dv d\omega = v' \mu' \alpha(z, v') \Sigma_s(z, v') \quad (1)$$

does lead to (K-7) and (K-11). (In an appendix at the close of this letter, we answer a question regarding the definition of α in Eq. (1).) To verify

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¹M. D. KOSTIN, *Nucl. Sci. Eng.* **19**, 252-254 (1964).

²A. M. WEINBERG and E. P. WIGNER, *The Physical Theory of Neutron Chain Reactors*, p. 413, Univ. of Chicago Press, Chicago, (1958).

³E. A. GUGGENHEIM, in *Encyclopedia of Physics* (S. Flügge, ed.), Vol. III/2, p. 3, Springer-Verlag, Berlin, (1959).

this statement, apply the operator $\int \mu v dv d\omega$ to (K-4) to obtain

$$\begin{aligned} \frac{\partial J(z, t)}{\partial t} = & - \frac{\partial}{\partial z} \int \mu^2 v^2 n(z, v, \underline{\omega}, t) dv d\omega - \\ & - \int \mu v^2 \Sigma(z, v) n(z, v, \underline{\omega}, t) dv d\omega + \\ & + \int \mu v v' \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) \times \\ & \times n(z, v', \underline{\omega}', t) dv' d\omega' dv d\omega. \quad (2) \end{aligned}$$

By use of the definition of $\alpha(z, v')$ from Eq. (1) we may write the final term as

$$\begin{aligned} & \int \mu' (v')^2 \alpha(z, v') \Sigma_s(z, v') n(z, v', \underline{\omega}', t) dv' d\omega' \\ & = \int \mu v^2 \alpha(z, v) \Sigma_s(z, v) n(z, v, \underline{\omega}, t) dv d\omega \end{aligned}$$

and obtain (K-7) and (K-11) unaltered. The definition (K-5) does not allow this transformation of the final term of Eq. (2).

(III) If the goal is to remove the one-velocity restriction from (K-1), then we feel that a better way is to apply the operator $\int \mu d\omega$ rather than $\int v \mu dv d\omega$ to (K-4). Furthermore, we feel that it is preferable to include on the right-hand side of (K-4) a source term $S(z, v, \underline{\omega}, t)$. This source may be caused by fissions, and we assume it to be isotropic in the sense that $\int \mu S d\omega \equiv 0$. In the time-independent case, the immediate result of applying $\int \mu d\omega$ to this modified (K-4) is

$$\begin{aligned} 0 = & -v \frac{\partial}{\partial z} \int \mu^2 n(z, v, \underline{\omega}) d\omega - v \Sigma(z, v) \int \mu n(z, v, \underline{\omega}) d\omega + \\ & + \int \mu v v' n(z, v', \underline{\omega}') \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) dv' d\omega' d\omega. \end{aligned}$$

With the definitions

$$J(z, v) = v \int \mu n(z, v, \underline{\omega}) d\omega,$$

$$N(z, v) = \int n(z, v, \underline{\omega}) d\omega,$$

$$\Sigma_s(z, v \rightarrow v) = \int \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega,$$

$$\int \mu \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega = \mu' \bar{\mu}_0(z, v' \rightarrow v) \Sigma_s(z, v' \rightarrow v), \quad (3)$$

$$\text{and } \bar{\mu}^2(z, v) = \int \mu^2 n(z, v, \underline{\omega}) d\omega / N(z, v),$$

this result becomes

$$\begin{aligned} 0 = & -v \frac{\partial}{\partial z} [\bar{\mu}^2(z, v) N(z, v)] - \Sigma(z, v) J(z, v) + \\ & + \int J(z, v') \bar{\mu}_0(z, v' \rightarrow v) \Sigma_s(z, v' \rightarrow v) dv'. \quad (4) \end{aligned}$$

(In an appendix at the close of this letter, we give concise answers to certain questions regarding the definition of $\bar{\mu}_0$ in Eq. (3).)

To obtain a "generalized Fick's law", we define

$$\Sigma_{tr}(z, v) = \Sigma(z, v) - \int J(z, v') \bar{\mu}_0(z, v' \rightarrow v) \times \\ \times \Sigma_s(z, v' \rightarrow v) dv' / J(z, v)$$

and write

$$J(z, v) = - \frac{v}{\Sigma_{tr}(z, v)} \frac{\partial}{\partial z} [\bar{\mu}^2(z, v) N(z, v)]. \quad (5)$$

Equation (5) does not have the simplicity of Fick's law for the reasons already listed in comment (I). In terms of the quantities defined above, Eq. (1) becomes

$$\int v \mu' \bar{\mu}_0(z, v' \rightarrow v) \Sigma_s(z, v' \rightarrow v) dv = v' \mu' \alpha(z, v') \Sigma_s(z, v') . \quad (6)$$

Consequently (K-11) may be derived from Eq. (4) by application of the operator $\int v dv$. Since information is lost in application of the operator $\int v dv$, Eq. (4) cannot be derived from (K-11). If we assume that scattering takes place without change in speed, then we may write

$$\int \mu \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega = \int \mu \delta(v' - v) \times \\ \times \Sigma_s(z, v', \underline{\omega}' \cdot \underline{\omega}) d\omega = \mu' \delta(v' - v) \bar{\mu}_0(z, v') \Sigma_s(z, v') .$$

It follows that for scattering without change in speed, (K-1), (K-2) and (K-3) are valid with J , Σ_{tr} , $\bar{\mu}^2$, N , Σ , $\bar{\mu}_0$ and Σ_s all functions of z and v . The v dependence is meaningful because of the v dependence of the fission source term. Here $\Sigma_{tr}(z, v)$ is an intensive property of the medium.

(IV) Instead of writing and using a 'generalized Fick's law' in the form of (K-11) or of Eq. (5), we believe that it is more direct to use the first velocity moment Eq. (2) or the first angular moment Eq. (4). Each of these first-moment equations is equivalent to a corresponding 'generalized Fick's law'. Usually a first-moment equation (or a 'generalized Fick's law') is not sufficient by itself for solution of problems in reactor analysis, so additional equations are needed. An appropriate zeroth-moment equation, analogous to (K-15), is frequently employed. Neither these moment equations nor their application in reactor analysis nor the idea of neutron pressure is new. Persiani⁴ derives these moments of the Boltzmann equation in three-dimensional form. He also gives application in reactor analysis and discusses the physical significance of these and of higher moments in terms of neutron pressure and other quantities. For a review and further application of methods of angular and velocity moments in

transport theory, see Ziering and Kahalas⁵ and Johnston⁶.

(V) The part of Ref. 1 from (K-1) through (K-3) and from (K-13) through (K-20), including Table I, is essentially independent of the remainder of the reference. The use of (K-11) in the two examples is clearly not essential. Table I, giving $\bar{\mu}^2(z)$ as a function of $z \Sigma_s$, is instructive.

(VI) Weinberg and Noderer⁷ and Amaldi⁸ present helpful and appropriate discussions of forms of Fick's law.

APPENDIX

Here we give concise answers to the following questions regarding the definitions of α and $\bar{\mu}_0$ in Eqs. (1) and (3): 1) are α and $\bar{\mu}_0$ independent of $\underline{\omega}'$ and $\underline{\omega}$? 2) Is the $\bar{\mu}_0$ of Eq. (3) the same as the more conventional $\bar{\mu}_0$ defined by

$$\frac{\int \underline{\omega}' \cdot \underline{\omega} \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega}{\int \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega} \\ = \frac{2\pi \int_{-1}^{+1} \mu_0 \Sigma_s(z, v' \rightarrow v, \mu_0) d\mu_0}{2\pi \int_{-1}^{+1} \Sigma_s(z, v' \rightarrow v, \mu_0) d\mu_0} \quad ?$$

Now by the Fenyő-Khekke theorem⁹

$$\int \mu \Sigma_s(z, v' \rightarrow v, \underline{\omega}' \cdot \underline{\omega}) d\omega \\ = 2\pi \mu' \int_{-1}^{+1} \mu_0 \Sigma_s(z, v' \rightarrow v, \mu_0) d\mu_0 .$$

It follows that the answers to the questions about $\bar{\mu}_0$ are affirmative. Since $\bar{\mu}_0$ is independent of $\underline{\omega}'$ and $\underline{\omega}$, the relation between α and $\bar{\mu}_0$ in Eq. (6) implies that α is independent of $\underline{\omega}'$ and $\underline{\omega}$.

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⁶T. W. JOHNSTON, *Phys. Rev.* **120**, 1103-1111 (1960).

⁷A. M. WEINBERG and L. C. NODERER, "Theory of Neutron Chain Reactions," Vol. I, pp. 5-25, CF-51-5-98 (AECD-3471), Oak Ridge National Laboratory (1951).

⁸E. AMALDI, in *Encyclopedia of Physics* (S. Flügge, ed.), Vol. XXXVIII/2, pp. 520-527, Springer-Verlag, Berlin, (1959).

⁹V. S. VLADIMIROV, "Mathematical Problems in the One-Velocity Theory of Particle Transport," AECL-1661, pp. 187 and 288, Atomic Energy of Canada Limited, Chalk River, Ontario (1963).

⁴P. J. PERSIANI, "Special Lectures on The Physical Foundation of Reactor Analysis," ANL-6227, Argonne National Laboratory, (1960).