

## Letters to the Editor

### Comment on Carlvik's Paper on Dancoff Correction in Lattices

This letter describes a simple FORTRAN routine for the function  $Ki_3(x)$  referred to by Carlvik.<sup>1</sup> It has at least the accuracy required.<sup>2</sup> Such a routine might not be conveniently available to many readers.

#### FUNCTION BIC3(X)

```
BIC3=(.52197980+X*(3.4150861+X*(2.8822975+X*.49523037)))/
1 (.66459895+X*(4.1979495+X*(2.7422822+X*.39516928)))/
2 (SORT(1.+X)*EXP(X))
RETURN
END
```

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<sup>1</sup>I. CARLVIK, *Nucl. Sci. Eng.*, 29, 325 (1967).

<sup>2</sup>I. GARGANTINI and T. PONENTALE, *Comm ACM*, 7, 12, 727 (1964).

### On the Connection Between Absorption and the Anisotropy of Scattering in Neutron Transport Theory

In a previous paper<sup>1</sup> the following rather remarkable theorem was proved. Consider one-speed neutron transport in a medium with a scattering law characterized by Legendre components  $p_1, p_2, \dots, p_n$ . In other words,

$$f(\mu_0) = \sum_{\ell=0}^N \frac{2\ell+1}{4\pi} p_\ell P_\ell(\mu_0), \quad (1)$$

where  $\mu_0$  is the scattering angle and  $P_\ell$  a Legendre polynomial. We denote by  $c$  the ratio of scattering-to-absorption in this medium. Then, the diffusion area  $\gamma_0^2$  can be expressed in a power series in  $(1-c)$  of the form

$$\gamma_0^2 = \sum_{m=1}^N \alpha_m (1-c)^m. \quad (2)$$

The theorem states that the  $\alpha_m$  are the same function only of  $p_1, p_2, \dots, p_m$  for any scattering law.

To state it in another way, consider  $\gamma_0^2$  for a medium described by  $p_1, \dots, p_n, p_{n+1} = p_{n+2} = p_{n+3} = \dots = 0$ . Consider another medium described by  $p_1, \dots, p_n, p_{n+1} \neq 0, p_{n+2} = p_{n+3} = \dots = 0$ . The theorem states that the first  $n$  terms of Eq. (2) are identical for the two media.

<sup>1</sup>E. INÖNÜ and A. İ. USSELLI, *Nucl. Sci. Eng.*, 23, 251 (1965).

In a subsequent work<sup>2</sup> the same theorem was shown to apply to the "effective source strength"<sup>3</sup> for the asymptotic solution. Thus, in fact, the theorem also applies to the asymptotic Green's function flux in an infinite medium.

The utility of this theorem is that it permits one to determine immediately the magnitude of the effect a given anisotropy will produce on the asymptotic flux. For example, if it is known that a problem can be described in sufficient accuracy by terms of order  $(1-c)^N$ , then all anisotropic components of the scattering  $p_j, j > N$  can be ignored.

The purpose of this letter is to point out that a similar simplification does not appear to hold for half-space problems. The proof is simply the production of a *Gegenbeispiel*. We consider the expression<sup>4</sup> for the extrapolated endpoint  $z_0(c)$ . For isotropic scattering it has the form

$$CZ_0(c) = 0.710446[1 - 0.0199(1-c)^2 + \dots] \quad (3a)$$

For linearly anisotropic scattering, we have<sup>4</sup>

$$CZ_0(c) = \frac{0.710446}{1-p_1}[1 - 0.0199(1-c)^2 + \dots] \quad (3b)$$

Consider the case  $c=1$ , with linearly anisotropic scattering. Then,

$$Z_0 = \frac{0.710446}{1-p_1} \quad (4)$$

If the theorem were true for this half-space problem, the value of  $z_0$  for  $c=1$  for quadratically anisotropic scattering would not depend upon  $p_2$ . However, in Case and Zweifel,<sup>4</sup> page 172,  $z_0$  is calculated for quadratically anisotropic scattering, and it is seen that it does indeed depend upon the value of  $p_2$ . This proves that the theorem does not apply to the Milne problem and presumably it does not apply to other half-space problems or to problems involving finite media.

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<sup>2</sup>A. İ. USSELLI, Thesis, Middle East Technical University (1967).

<sup>3</sup>A. M. WEINBERG and E. P. WIGNER, *The Physical Theory of Neutron Chain Reactors*, The University of Chicago Press, Chicago, Illinois (1958).

<sup>4</sup>K. M. CASE and P. F. ZWEIFEL, *Linear Transport Theory*, Addison Wesley, Reading, Massachusetts (1967).