

TABLE I  
The  $H$ -Function and its Analytic Approximation

$c$	0	0.2	0.4	0.6	0.8	1.0
$\mu$	$H_{22}^{(1)} = \left[ 1 - \frac{c\alpha}{2}\mu - (1-\alpha\mu)\frac{c}{2}\mu \ln \frac{1+\mu}{\mu} \right]^{-1}$					
0	1	1	1	1	1	1
0.2	1	1.03871	1.0848	1.1427	1.2233	1.444
0.4	1	1.05523	1.1241	1.2155	1.3540	1.819
0.6	1	1.06564	1.1498	1.2657	1.4511	2.179
0.8	1	1.07297	1.1683	1.3031	1.5276	2.534
1.0	1	1.07844	1.1823	1.3323	1.5899	2.885
$H$ -function						
0	1	1	1	1	1	1
0.2	1	1.03892	1.08577	1.14517	1.2286	1.4503
0.4	1	1.05546	1.12516	1.21861	1.3611	1.8293
0.6	1	1.06587	1.15087	1.26893	1.4590	2.1941
0.8	1	1.07319	1.16935	1.30631	1.5358	2.5527
1.0	1	1.07864	1.18337	1.33541	1.5982	2.9078

TABLE II

First Moment of the  $H$ -Function;  $\alpha_1^{[*]} = \int_0^1 H(\mu)\mu d\mu$ ;

$$\alpha_1 \simeq \bar{\alpha}_1 = \int_0^1 \frac{1}{2} (H_{21}^{(1)} + H_{11}^{(1)}) \mu d\mu$$

$c$	0	0.1	0.2	0.3	0.4	0.5
$\alpha_1$	0.5	0.515609	0.533154	0.553123	0.576210	0.603495
$\bar{\alpha}_1$	0.5	0.51553	0.53316	0.55314	0.57624	0.60352
$c$		0.6	0.7	0.8	0.9	1.0
$\alpha_1$		0.636636	0.678674	0.735808	0.825318	1.154701
$\bar{\alpha}_1$		0.63665	0.67862	0.7355	0.8242	1.145

<sup>[\*]</sup>Deduced from Tables XI and XXXIII of Chandrasekhar<sup>3</sup> which are accurate to within 0.005%.

TABLE III  
(b)  $1 - R(\mu_0 = 1)$

$c$	0	0.2	0.4	0.6	0.8	1.0
Rafalski	1	0.96306	0.9076	0.8222	0.6716	0
Pomraning	1	0.96391	0.9137	0.8398	0.7078	0
$H \approx H_{22}^{(1)}$	1	0.96459	0.9158	0.8426	0.7110	0
Chandrasekhar <sup>a</sup>	1	0.96476	0.9166	0.8446	0.7147	0

<sup>a</sup>Deduced from Tables XI and XXXIII of Chandrasekhar<sup>3</sup> which are accurate to within 0.005%.

TABLE IV

$$1 - \bar{A}_{is} = 2\alpha_1 \sqrt{1-c}$$

$c$	0	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9	1
Pomraning	1	0.980	0.957	0.929	0.896	0.856	0.808	0.746	0.660	0.523	0
$\alpha_1 \approx \bar{\alpha}_1$	1	0.97814	0.95375	0.92558	0.89270	0.85351	0.80494	0.7434	0.6579	0.5213	0
Chandrasekhar <sup>a</sup>	1	0.97830	0.95374	0.92555	0.89266	0.85347	0.80529	0.74345	0.65813	0.52198	0

<sup>a</sup>Deduced from Tables XI and XXXIII of Chandrasekhar<sup>3</sup> which are accurate to within 0.005%.

$H \approx H_{11}^{(1)}$  or  $H_{21}^{(1)}$ . For a simple expression for the first moment of the  $H$ -function, we found it best to approximate  $H$  by  $H \approx \frac{1}{2}[H_{11}^{(1)} + H_{21}^{(1)}]$ . This choice leads to

$$\alpha_1 \simeq \bar{\alpha}_1 = \frac{1}{2} + \left(\frac{\alpha}{3} + \frac{c}{6}\right) \left(\ln 2 - \frac{1}{4}\right) + \frac{c\alpha}{12} - \frac{c\alpha^2}{10} \left(\ln 2 - \frac{11}{12}\right). \quad (15)$$

This result is compared in Table II with the first moments of the  $H$ -function  $\alpha_1$  given by Table XXXIII of Chandrasekhar<sup>3</sup>.

Equations (9) through (15) when applied to Eqs. (1) and (2) constitute the present approximations for the albedos. The approximations of Eqs. (1), (14) and (2), (15) for albedo defects  $1-R(1)$  and  $1-\bar{A}_{is}$  are compared in Tables III and IV with the corresponding results of Rafalski and Pomraning, and with values computed from Tables XI and XXXIII of Chandrasekhar<sup>3</sup> as a reference.

This work was done under the supervision of G. Birkhoff. Further work on the above method of approximation has been done and an extension is in progress.

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March 18, 1966

### Cavity and Vacuum Boundary Conditions for More-dimensional $P_N$ Approximations

In the June 1966 issue of *Nuclear Science and Engineering*, Davis<sup>1</sup> has derived vacuum boundary conditions for the neutron flux and its adjoint in a  $P_N$  approximation ( $N$  even and odd) by variational methods.

Equivalent boundary conditions have been obtained by neutron balance considerations<sup>2</sup>. A natural approach is to start with cavity boundary conditions for exposed surfaces. The number of neutrons that enter a surface element  $dS = \mathbf{n} \cdot d\mathbf{S}$  per  $s$  in a solid angle  $d\Omega$  around  $\Omega$  is determined by neutrons that leave the surface element  $dS' = \mathbf{n}' \cdot dS'$  in the solid angle  $-(\Omega \cdot \mathbf{n})dS/R^2$  around  $\Omega$

$$-\Phi(\mathbf{r}, \Omega)(\Omega \cdot \mathbf{n})d\Omega dS = \Phi(\mathbf{r}', \Omega)(\Omega \cdot \mathbf{n}') \left[ -\frac{(\Omega \cdot \mathbf{n})}{R^2} \right] dS \cdot dS' \quad (1)$$

Hence,  $d\Omega$  is related to  $dS'$  by

$$d\Omega = (\Omega \cdot \mathbf{n}') \cdot dS'/R^2 \quad (2)$$

<sup>1</sup>J. A. DAVIS, "Variational Vacuum Boundary Conditions for a  $P_N$ -Approximation," *Nucl. Sci. Eng.*, 25, 2, 189 (1966).

<sup>2</sup>D. EMENDÖRFER, "Randbedingungen für den Neutronenfluss im endlichen Zylinder nach der  $P_N$ -Approximation der Transportgleichung," *Nukleonik*, 5, 74 (1963).